Analytical and Numerical Studies of Ion Beam Plasma Interaction for <u>Heavy Ion Driven</u> Inertial Fusion

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Talk Outline

Basics

Why do Heavy Ion Fusion? What do we have to do? A bit about the accelerator and the beams

Program

Past Accomplishments

Present Program

Future Plans

Physics of Ion Beam Plasma Interaction

Degree of charge and current neutralization Self electric and magnetic fields



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Introduction to Heavy Ion Fusion

 Inertial fusion is based on H-bomb design

Who built the h-bomb? Debate revives
By WILLIAM J. BROAD, April 24, 2001

The New York Eimes





Bomb Basics

PRIMARY



Conventional explosives compress plutonium in the primary, creating a critical mass in which atoms begin to split apart and release nuclear energy.



The radiation vaporizes the lining of the casing and radiates back toward the secondary, compressing it and heating it to fusion temperature.

SECONDARY



Radiation from the primary flows down the length of the bomb casing ahead of the primary blast.



Heavy Ion Fusion Concept







How IFE Targets Work















The Heavy Ion Fusion Virtual National Laboratory

400-MJ IFE Capsule Dimensions



Heavy Ion Fusion Requirements

Target requirements3 - 7 MJ $x \sim 10 \text{ ns}$ $\Rightarrow \sim 500 \text{ Terawatts}$ (hand grenade) (Forty times the averaged world-wide electric power consumption)

Ion Range: 0.02 - 0.2 g/cm² \Rightarrow 1- 10 GeV







Peaceful Power -- An Artist's Conception of a HIF Power Plant





Why Heavy Ion Fusion?



Advantages of Heavy Ion Fusion Approach

• Driver is separate from the fusion chamber

⇒ accelerator not exposed to fusion environment
⇒ can protect 1st wall

- High electrical efficiency
- Takes advantage of decades of worldwide investment in accelerators & defense-funded target design



Heavy Ion Accelerators are a Good Choice for a Fusion Driver

High Energy Physics accelerators already have:

Long life High pulse repetition rates High efficiency (~ 30%) Present systems comparable to requirements in: complexity cost ion energy

But: much higher charge-per-unit-length is needed. induction is not used in HEP accelerators



The First Wall is Protected by Neutronthick Molten Salt (FLiBe)



Fusion Pocket Formed by Liquid Jets

Flibe (Li2BeF4), is a molten salt with twice the density of water, and roughly the same viscosity, desired for making high-pressure steam for driving turbines.

Flibe absorbs neutrons.





Why Not Laser Drivers?

Lasers:

Easy to focus Much more money in program Easier development path

but:

Problem protecting final optics
Problem protecting first wall
Target gets damaged (hot)AllLow repetition rate (a few/day)Low electrical efficiency (a few - 10%)

Glass





Features of the Accelerator

3 - 7 MJ / 1- 10 GeV ions \Rightarrow ~ 10 ¹⁶ ions / 100 beams

⇒ 1-2 kA / beam =>Space-charge-dominated beams (non-neutral plasma)

For accelerator design:

Focusability ⇒ Multiple beams

Efficiency \Rightarrow **Induction acceleration**

Stability \Rightarrow Linear accelerator





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Present Experiments

Present experiments – higher current (larger radius) similar to the driver at low energy ⇒ high space charge potential.

0.1 - 0.5 A, 0.4 - 1 MeV, space charge potential ~ few kV

High Current Experiment (HCX) – how much beam can be transported?
Neutralized Transport Experiment (NTX) – neutralization by plasmas after the final focus



The Present Experimental Program: HCX and NTX experiments



Many, but not all, issues experimentally explored

Very complete experimental understanding



The High Current Experiment HCX)

Issues to be resolved:

How much beam can be transported? Dynamic aperture (usable aperture set by dynamics) Halo production Effect of desorbed gas on tail Electron production & orbits (magnetic transport) Mismatch, misalignment

Parameters:

K⁺ or Cs⁺ ~ 0.2 - 0.5 Amp 1 - 1.7 MeV 4.5 - 7 μs

Quadrupoles:

- 10 (30-40 later) electrostatic
 - 4 pulsed normal magnetic
 - a few superconducting



High Current Experiment (HCX) operation since January, 2002



Transport

- aperture limits
- electrons
- gas effects
- halo formation
- steering



Chamber Transport



Neutralization competes with stripping in the target chamber





Neutralized Transport Experiment



400 kV injector

Final-focusing optical system

FY02: characterize ion & plasma sources FY03: study beam aberrations FY04: complete initial neutralization experiments



Beam Focusing with Plasma Plug (operation since September, 2002)



Pulsed arc

plasma source



Without Plasma

With Plasma



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Scintillating glass

Drift tube

Beam simulations and theory span a variety of



acceleration in 3-D structure



two-stream instability



chamber propagation

electrons during neutralization

Integrated Beam eXperiment Short-Pulse Source-to-Focus Experiment





Heavy Ion Fusion Summary

- Because of the separation of accelerator from fusion chamber, the heavy ion fusion concept is able to protect the 1st wall, and the driver, from fusion products.
- HIF benefits from large U.S. and worldwide investments in accelerators and defense-funded target research.
- Past and present experiments have demonstrated production, acceleration, compression, and focusing of beams at lower energy and current.
- The Integrated Beam Experiment would be a proof-of-principle experiment with a rich physics mission: to explore longitudinal physics and test integrated source-to-focus transport.





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Means of Investigation

Nonlinear theory
 Numerical simulation:

 Particle in cell
 Fluid

 Neutralization experiment (VNL)



Results of 2D PIC-MC Code

Beam propagation in the y-direction,
beam length 7.5 c/ω_p;
beam radius 1.5 c/ω_p;
beam density equal to the half of the plasma density;
beam velocity c/2.

•Shown are electron density and the current.









beam length $30. c/\omega_p;$ beam radius $0.5 c/\omega_p;$ •beam density is 5 of plasma density; •beam velocity 0.5c.



The Heavy lon

System of Equations

$$\frac{\partial n_e}{\partial t} + \nabla \bullet \left(n_e \vec{V_e} \right) = 0,$$

$$\frac{\partial \vec{p}_e}{\partial t} + (\vec{V}_e \bullet \nabla) \vec{p}_e = -\frac{e}{m} (\vec{E} + \frac{1}{c} \vec{V}_e \times \vec{B}),$$

$$\nabla \times \vec{B} = \frac{4\pi e}{c} \left(Z_b n_b V_{bz} - n_e V_{ez} \right) + \frac{1}{c} \frac{\partial \vec{E}}{\partial t},$$

 $\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}.$ The Heavy Ion Fusion Virtual National Laboratory



Steady- State Results (current flow)

Beam propagates in the y-direction, beam half length $l_b=15 c/\omega_p$; beam radius $r_b=1.5 c/\omega_p$; beam density n_b is equal to the background plasma density n_p ; beam velocity $V_b=c/2$.

Shown are the normalized electron density n_e/n_p and the vector fields for the current.

FOR MORE INFO...

http://hifnews.lbl.gov/hifweb08.html





Steady- State Results (electric field)

Beam propagates in the y-direction, $l_b=15 c/\omega_p$; $r_b=1.5c/\omega_p$; $n_b=n_p$; $V_b=c/2$.

Shown are the normalized electron density n_e/n_p and the vector fields for the electric force on electrons.

FOR MORE INFO...

http://hifnews.lbl.gov/hifweb08.html





Steady- State Results



normalized electron current $j_v/(ecn_p)$

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http://www.trilobites.com



Analytic theory of chamber transport: excitation of plasma waves by beam depends on bunch length



 $\beta_b = 0.5$, $I_b/r_b = 10$, $n_b/n_p = 0.5$ a) $I_b = 2V_b/\omega_p$, b) $I_b = 6V_b/\omega_p$, c) $I_b = 20V_b/\omega_p$. Red line: ion beam size, brown lines: electron trajectory in beam frame



Nonlinear Theory If you would like to experience delicious taste, bite in small pieces.

• Important issues:

- Finite length of the beam pulse,
- Arbitrary value of n_b/n_p ($n_b > > n_p$),
- 2D.
- Approximations:
 - Fluid approach,
 - Conservation of generalized vorticity,
 - Long dense beams $l_b >> r_b$, V_b/ω_p .

• Exact analytical solution.



Conservation of Generalized Vorticity

$$\begin{split} \vec{\Omega}_{e} &= \nabla \times \vec{p}_{e} - e\vec{B}/c, \\ \frac{\partial \vec{\Omega}_{e}}{\partial t} - \nabla \times (\vec{V}_{e} \times \vec{\Omega}_{e}) = 0, \\ & \oint \vec{\Omega}_{e} \cdot d\vec{S} = \oint (\vec{p}_{e} - \frac{e}{c}\vec{A}) \cdot \delta\vec{r} = const., \\ & \vec{\Omega}_{e} = 0 \Longrightarrow \nabla \times \vec{p}_{e} = e\vec{B}/c. \end{split}$$

FOR MORE INFO... I. Kaganovich, et.al, Physics of Plasmas 8, 4180 (2001).

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Approximate System of Equations

$$\begin{split} l_b >> r_b, \qquad & (V_{ey} - V_b) \frac{\partial n_e}{\partial y} + \frac{1}{r} \frac{\partial}{\partial r} \left(rn_e V_{er} \right) = 0, \\ \nabla \times \vec{\mathbf{p}}_e &= \frac{e}{c} \vec{\mathbf{B}}, \qquad \frac{1}{r} \frac{\partial}{\partial r} E_r = 4\pi e (Z_b n_b + n_p - n_e), \\ K_e &= m_e \vec{V_e}^2 / 2, \qquad & (V_{ey} - V_b) \frac{\partial p_{er}}{\partial y} + \frac{\partial K_e}{\partial r} = -eE_r, \\ & -\frac{1}{r} \frac{\partial}{\partial r} r \frac{\partial}{\partial r} p_{ey} = \frac{4\pi e}{c} \left(Z_b n_b V_{by} - n_e V_{ey} \right) \end{split}$$



Simplified Code

• Approximation of long beams:

- Beam length is much longer than beam radius;
- Therefore, beam can be described by a number of weakly interacting slices.
- The electric field is found from radial Poisson's equation.
- As a result of the simplification the second code is hundreds times faster than the first one and can be used for most cases, while the first code provides benchmarking for the second.



Comparison of Theory and Simulation: Electron Density

Electron density Left – PIC, Right - fluid

 $l_b = 1c/\omega_p, r_b = 0.1c/\omega_p$ $n_b = 0.5n_p, V_b = 0.5c$

Brown lines: electron trajectory in the beam frame. Red line: ion beam size.

I. Kaganovich et.al,



http://pacwebserver.fnal.gov/papers/Tuesday/PM_Poster/TPPH317.pdf



Comparison of Theory and Simulation: Electron Density

Key parameter $\omega_p l_b / V_b$, Quasineutrality $l_b >> V_b / \omega_p$.









Comparison of Theory and Simulation: Current Neutralization

Key parameter $\omega_p r_b/c$, Current neutralization $r_b >> c/\omega_p$.









Comparison of Theory and Simulation: Magnetic Field

Key parameter $\omega_p r_b/c$, Magnetic field neutralization $r_b >> c/\omega_p$.



FOR MORE INFO... I. Kaganovich, et.al, Physics of Plasmas 8, 4180 (2001).



Results and Conclusions

- Developed a nonlinear fluid theory for the quasi-steady-state propagation of an intense ion beam pulse in a background plasma under the assumption of a long beams $l_b >> r_b$.
 - The analytical formulas can provide an important benchmark for numerical codes.
 - The analytical solutions form the basis of a hybrid semi-analytical approach, used for calculations of beam propagation in the target chamber.
- The simulations of current and charge neutralization performed for conditions relevant to heavy ion fusion showed:
 - very good charge neutralization: key parameter $\omega_p l_b/V_b$,
 - very good current neutralization: key parameter $\omega_p r_b/c$.
- Plasma wave breaking heats electrons $n_b > n_{p^*}$



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beam velocity c/2.

•Shown is the electron current.





Plasma Wave Breaking Heats the Electrons when $n_b > n_p$



Electron phase space shown for $l_b=30V_b/\omega_p$; $n_b=2n_p$. Times after entering the plasma plug are: (a) $113/\omega_p$, and (b) $245/\omega_p$.



Fluid approximation is good for

 $n_h \leq n_n$



Electron current and phase space $l_b=15c/\omega_p$; $n_b=n_p$; $V_b=c/2$.



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Beam propagation in the y-direction,
beam half length 30 c/ω_p;
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beam density is equal to 5 of the plasma density;
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